# Physics 603 Classical Mechanics - Spring 2020 Minkowski space and relativistic dynamics I

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## 1 3+1 Minkowski space

Vectors in Minkowski space have four components, which are indicated by a greek superscript  $v^{\mu}$  with  $\mu = 0, 1, 2, 3$ . A latin index is often used to distinguish the spatial part of the vector  $v^k$  with k = 1, 2, 3. The four-vector

$$x^{\mu} = (ct, \mathbf{x}),\tag{1.1}$$

represents an **event** which specifies the **position** where something happens and the **time** when it happens. Another example is the energy-momentum four-vector

$$p^{\mu} = (E/c, \mathbf{p}) = m_0 c(\gamma, \gamma \beta). \tag{1.2}$$

In general, four-vectors in Minkowski space must transform properly under Lorentz Transformations (boosts and rotations).

#### **1.1 Scalar product**

The inner product g of any four-vectors is an invariant

$$g(v^{\mu}, u^{\nu}) = g(v'^{\mu}, u'^{\nu}) = s.$$
(1.3)

For two events  $x^{\mu}, y^{\mu}$  the scalar product is given by

$$g(x^{\mu}, y^{\mu}) = (ct_x)(ct_y) - \mathbf{x} \cdot \mathbf{y} = x^0 y^0 - x^k y^k,$$
(1.4)

where the sum over k = 1, 2, 3 is implied. For the four-momentum we notice that the inner product with itself is proportional to the squared rest mass

$$g(p^{\mu}, p^{\mu}) = \frac{E^2}{c^2} - \mathbf{p}^2 = m_0^2 c^2.$$
(1.5)

The scalar product of the difference between two events dx with itself returns the **invariant** distance

$$g(dx^{\mu}, dx^{\mu}) = ds^{2} = d(ct)^{2} - d\mathbf{r}^{2}.$$
(1.6)

Two events can be classified according to their invariant distance as

1. Light-like:  $ds^2 = 0$  2. Space-like:  $ds^2 < 0$  3. Time-like:  $ds^2 > 0$ The eigentime, or proper time, for a time-like separation (distance) is defined as

$$d\tau = \sqrt{ds^2}/c = \gamma^{-1}dt. \tag{1.7}$$

#### 1.2 Four-forms

A four-form (or covariant vector) is a functional U that maps a Minkowski four-vector (or contravariant vector) into scalars

$$U(x^{\mu}) = s. \tag{1.8}$$

Like vectors, forms have four components indicated by a lower index  $u_{\mu}$  such that

$$U(v^{\mu}) = u_{\mu}v^{\mu}.$$
 (1.9)

#### 1.3 Tensors

Tensors are objects with more than one upper or lower index

$$T^{\mu\nu\dots}_{\alpha\beta\dots},\tag{1.10}$$

these can be constructed out of vectors and forms. The Lorentz matrix itself represents a tensor which takes a vector in a frame S and returns a vector in a different frame S', thus

$$S': \qquad x'^{\mu} = \Lambda^{\mu}_{\nu} x^{\nu}. \tag{1.11}$$

For example, a boost in the z direction is given by

$$\Lambda^{\mu}_{\nu} = \begin{pmatrix} \gamma & 0 & 0 & -\gamma\beta \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ -\gamma\beta & 0 & 0 & \gamma \end{pmatrix}.$$
 (1.12)

The metric tensor g

$$g_{\mu\nu} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$$
(1.13)

allows us to write the scalar product of two vectors as

$$g(x^{\mu}, y^{\nu}) = g_{\mu\nu} x^{\mu} y^{\nu}. \tag{1.14}$$

On the other hand, the action of the metric tensor on a vector returns a form, thus lowers its index

$$g_{\mu\nu}x^{\nu} = x_{\mu} = (ct, -\mathbf{x}).$$
 (1.15)

### 2 Relativistic formulation

Upon having defined the proper objects of the Minkowski (scalars, vectors, forms, and tensors) now it is possible to rewrite the laws of physics in a relativistic formulation. The non-relativistic coordinates are replaced by events  $x^{\mu}$ , with the **four-velocity** given by the derivative with respect to the eigentime

$$\frac{dx^{\mu}}{d\tau} = \gamma \frac{dx^{\mu}}{dt} = (\gamma c, \mathbf{u}), \qquad (2.1)$$

such that

$$m_0 u^{\mu} = (\gamma m_0 c, \gamma m_0 \mathbf{u}) = (E/c, \mathbf{p}) = p^{\mu}.$$
 (2.2)

Second Newton's law can then be written in the relativistic formulation as

$$K^{\mu} = \frac{dp^{\mu}}{d\tau} = \gamma \frac{dp^{\mu}}{dt}, \qquad (2.3)$$

where  $K^{\mu}$  represents the **four-force**. This becomes even more evident by defining the **pseudo-force**  $F^{\mu} = K^{\mu}/\gamma$ 

$$F^{\mu} = \frac{dp^{\mu}}{dt} = \left(\frac{dp^{0}}{dt}, \mathbf{F}\right) = \left(\frac{1}{c}\frac{dE}{dt}, \mathbf{F}\right), \qquad (2.4)$$

where we can identify the power of the system in the zero component of the pseudo-force. In fact, the scalar product between the pseudo-force and the four-velocity returns the expression for the power in terms of the spatial velocity

$$\frac{dp^{\mu}}{dt}u_{\mu} = \gamma \left(\frac{dE}{dt} - \mathbf{F} \cdot \mathbf{u}\right) = \frac{m_0}{2}\frac{d}{dt}u^{\mu}u_{\mu} = \frac{m_0}{2}\frac{d}{dt}c^2 = 0 \quad \Rightarrow \quad \frac{dE}{dt} = \mathbf{F} \cdot \mathbf{u}.$$
 (2.5)