

$H(\vec{q}, \vec{p}, t)$ sometimes = E and sometimes conserved.

1) $\frac{\partial H}{\partial q_i} = -\dot{p}_i$ 2) $\frac{\partial H}{\partial p_i} = \dot{q}_i$ HEM

From variation principle

$\int_{t_1}^{t_2} \sum_i p_i \dot{q}_i - H(\vec{q}, \vec{p}, t) dt$ must be stationary ($\delta S = 0$)

along \forall path in phase space with

$\vec{q}(t_1), \vec{p}(t_1), \vec{q}(t_2), \vec{p}(t_2)$ fixed

q_1, \dots, q_k
 p_1, \dots, p_k } $2k$

We can describe system in phase space

1) shows that if there is no q_i in $H \Rightarrow p_i = \text{const}$
 \Rightarrow we can integrate the expression. If H doesn't depend on all q_i 's \Rightarrow the problem becomes trivial. It can happen when problem is trivial or if we choose the generalised coord-s well.

We must be able to change variables from one to another set. We must be able to replace

$\boxed{\vec{q}, \vec{p} \rightarrow \vec{Q}, \vec{P}}$ (\vec{q}, \vec{p}, t) \Rightarrow $H \rightarrow K(\vec{Q}, \vec{P}, t)$

K can be = to H . Require $\frac{\partial K}{\partial Q_i} = -\dot{P}_i, \frac{\partial K}{\partial P_i} = \dot{Q}_i$

If $Q_i = Q_i(\vec{q}, t)$ only \rightarrow point transformation in configuration space
 transf.-n is canonical means that also

$\delta \int_{t_1}^{t_2} \sum_i p_i \dot{Q}_i - K(\vec{Q}, \vec{P}, t) dt = 0$ (it also has to be stationary)

if we can add something to \int from which the δS automatically becomes 0 $\Rightarrow \delta S$ is still 0

$$\delta \int_{t_1}^{t_2} \left\{ \sum P_i \dot{Q}_i - K(\vec{Q}, \vec{P}, t) + \frac{dF}{dt} \right\} dt = 0.$$

$$\int = F(t_2) - F(t_1) \quad \leftarrow \quad F \text{ is generating function}$$

F is an arbitrary f.-n

which is const at t_2 and t_1
(beginning + end of path)

δS will not also change if we * it with λ .

We call it scale transformation

$$\delta \int \lambda (\sum p_i \dot{q}_i - H(\vec{q}, \vec{p}, t)) dt = 0.$$

Example 1

$$P_i = \gamma p_i, \quad Q_i = \nu q_i, \quad K = \gamma \nu H$$

$$\frac{\partial K}{\partial Q_i} = \frac{1}{\nu} \gamma \nu \frac{\partial H}{\partial q_i} = \gamma p_i = -\dot{P}_i$$

and the same for the second eq.-n.:

$$\frac{\partial K}{\partial P_i} = \frac{1}{\gamma} \frac{\partial K}{\partial p_i} = \frac{1}{\gamma} \gamma \nu \frac{\partial H}{\partial p_i} = \nu \dot{q}_i = \dot{Q}_i$$

$$\begin{aligned} \int \sum_i P_i \dot{Q}_i - K(\vec{Q}, \vec{P}, t) dt &= \int \sum \gamma p_i \nu \dot{q}_i - \gamma \nu H(\vec{q}, \vec{p}, t) dt \\ &= \int \lambda [\sum p_i \dot{q}_i - H] dt \end{aligned}$$

Assume we've chosen our coordinates such, that $\lambda = 1$.
For the new variables we should have following cond.:

$$\boxed{\sum p_i \dot{q}_i - H(\vec{q}, \vec{p}, t) = \sum P_i \dot{Q}_i - K(\vec{Q}, \vec{P}, t) + \frac{dF}{dt}} \quad (*)$$

if we can prove this, then we'll have canonical transform.-n

Example type 1 gen. f.-n $H_1(\vec{q}, \vec{Q}, t)$.

$$\text{Example: } F = \sum_i q_i Q_i \quad \frac{dF}{dt} = \sum_i \dot{q}_i Q_i + q_i \dot{Q}_i$$

In (*) neither K nor H depend on \dot{q}_i or \dot{Q}_i

$$\sum_i p_i \dot{q}_i - \sum_i \dot{q}_i Q_i = \sum_i P_i \dot{Q}_i + H(\vec{q}, \vec{P}, t) - K(\vec{Q}, \vec{P}, t) + \sum_i q_i \dot{Q}_i$$

nothing in right side depends on \dot{q}_i

$$\Rightarrow p_i - Q_i = 0$$

$$\Rightarrow \boxed{Q_i = p_i}$$

if we do same for $\dot{Q}_i \Rightarrow \sum_i (P_i + q_i) \dot{Q}_i \Rightarrow \boxed{P_i = -q_i}$

$$\Rightarrow \underline{H = K}$$

Example 3

Freely falling marble in 1-dimension.

$$H = \frac{p^2}{2m} + mgq \quad (q - \text{height})$$

$$Q = p, \quad P = -q$$

$$\Rightarrow K = \frac{Q^2}{2m} - mgP$$

$$\frac{\partial K}{\partial P} = \dot{Q} = -mg$$

↓
 $\dot{p} = -mg$

$$\frac{\partial K}{\partial Q} = -\dot{P} = \frac{Q}{m}$$

↓
 $\dot{q} = \frac{p}{m}$

Example 7 general case.

$$F_1(\vec{q}, \vec{Q}, t)$$

$$\frac{dF_1}{dt} = \sum_i \left(\frac{\partial F_1}{\partial q_i} \dot{q}_i + \frac{\partial F_1}{\partial Q_i} \dot{Q}_i \right) + \frac{\partial F_1}{\partial t}$$

\Rightarrow Again doing same steps as in example 2), from (*) we get;

$$\sum (p_i - \frac{\partial F_1}{\partial q_i}) \dot{q}_i \Rightarrow \boxed{p_i = \frac{\partial F_1}{\partial q_i}} = f_i(\vec{q}, \vec{Q}, t) \rightarrow \text{solve for } \vec{Q}(\vec{q}, \vec{p}, t)$$

$$\text{and } \sum (p_i + \frac{\partial F_1}{\partial Q_i}) \dot{Q}_i \Rightarrow \boxed{p_i = -\frac{\partial F_1}{\partial Q_i}} = g_i(\vec{Q}, \vec{q}, t) \rightarrow p_i(\vec{q}, \vec{p}, t)$$

$\Rightarrow K = H + \frac{\partial F_1}{\partial t}$, the transformation is canonical

These steps we do if we're given F .

But if we know Q, P , we need to find F ?

\Rightarrow we have $\frac{\partial F_1}{\partial q_i} = p_i$ and $p_i = -\frac{\partial F_1}{\partial Q_i}$ which we need to integrate to find F_1 .

these we can integrate if they fulfill condition

$$\frac{\partial^2 F_1}{\partial q_i \partial Q_i} = \frac{\partial^2 F_1}{\partial Q_i \partial q_i} \quad \text{or} \quad \frac{\partial}{\partial q_i} (-p_j) = \frac{\partial}{\partial Q_i} (p_i)$$

Example 5 Harmonic oscillator: $q, p = m\dot{q}$

$$p = m\omega A \cos \omega t$$

$$H = \frac{p^2}{2m} + \frac{1}{2} k q^2 = \frac{p^2}{2m} + \frac{1}{2} m \omega^2 q^2$$

$$q = A \sin \omega t$$

We want to choose Q such that $H = H(Q)$ and $P = \text{const}$

$$\omega^2 = \frac{k}{m}$$