

Graduate Quantum Mechanics - Problem Set 10

Problem 1)

Using the matrix elements of the operator \mathbf{L}_x in the subspace for $l = 1$ (see Shankar p. 327-328, in particular the 3rd block in the matrix 12.5.23), show that the matrix for arbitrary rotations around the x-axis (in the basis consisting of eigenstates $|l=1, m\rangle$ ($m = +1, 0, -1$) of \mathbf{L}_z) is given by

$$D_{mm'}(\theta) = \exp\left(-i\theta \frac{\mathbf{L}_x}{\hbar}\right) = \begin{pmatrix} \frac{1}{2}\cos\theta + \frac{1}{2} & -\frac{i}{\sqrt{2}}\sin\theta & \frac{1}{2}\cos\theta - \frac{1}{2} \\ -\frac{i}{\sqrt{2}}\sin\theta & \cos\theta & -\frac{i}{\sqrt{2}}\sin\theta \\ \frac{1}{2}\cos\theta - \frac{1}{2} & -\frac{i}{\sqrt{2}}\sin\theta & \frac{1}{2}\cos\theta + \frac{1}{2} \end{pmatrix}$$

Show that applying this matrix for the case $\theta=180^\circ$ on the eigenfunction $|l=1, m=1\rangle$ of \mathbf{L}_z gives the same result as rotating explicitly the function $Y_1^1(\theta, \varphi) = -\sqrt{\frac{3}{8\pi}} \sin\theta e^{i\varphi}$ by 180° around the x-axis.

Hint: Rotating by 180° around the x-axis is the same as replacing θ by $180^\circ - \theta$ and φ by $-\varphi$. Also, note that $Y_1^{-1}(\theta, \varphi) = \sqrt{\frac{3}{8\pi}} \sin\theta e^{-i\varphi}$.

Problem 2)

A particle in 3 dimensions is described by a wave function $\Psi(x, y, z)$. We can make a variable substitution to spherical variables $(x, y, z) \rightarrow (r, \theta, \varphi)$ by defining an alternative wave function in terms of the new variables which describes the exact same state:

$$\Phi(r, \theta, \varphi) = \Psi(x, y, z) = \Psi(r \sin\theta \cos\varphi, r \sin\theta \sin\varphi, r \cos\theta).$$

Using the definition of \mathbf{L}_x in position (x, y, z) space, show that its representation in terms of these new variables is

$$\mathbf{L}_x = -\frac{\hbar}{i} \left(\sin\varphi \frac{\partial}{\partial\theta} + \cos\varphi \cot\theta \frac{\partial}{\partial\varphi} \right).$$

Hint: Simply apply this form of \mathbf{L}_x to $\Phi(r, \theta, \varphi)$ and show that you get the same answer as applying \mathbf{L}_x in Cartesian coordinates to $\Psi(x, y, z)$.

Problem 3)

Using the result for \mathbf{L}_x above and the similar result for \mathbf{L}_y (see Exercise 12.5.8 p. 334 in Shankar – you don't have to prove that):

$$\mathbf{L}_y = \frac{\hbar}{i} \left(\cos\varphi \frac{\partial}{\partial\theta} - \sin\varphi \cot\theta \frac{\partial}{\partial\varphi} \right)$$

show that Eq. 12.5.27 is true, i.e.

$$\mathbf{L}_{\pm} = \pm\hbar e^{\pm i\varphi} \left(\frac{\partial}{\partial\theta} \pm i \cot\theta \frac{\partial}{\partial\varphi} \right)$$

for at least one of the two (you can then assume the other one is correct, as well).

Problem 4)

Using the results from Problem 3 and the fact that

$$\mathbf{L}_- \mathbf{L}_+ = \mathbf{L}_x^2 + \mathbf{L}_y^2 - \hbar \mathbf{L}_z$$

calculate the form of the total squared angular momentum operator $\mathbf{L}^2 = \mathbf{L}_x^2 + \mathbf{L}_y^2 + \mathbf{L}_z^2$ in spherical coordinates and show that the result agrees with Shankar's Eq. 12.5.36, p. 335