

QCD MADE SIMPLE

Frank Wilczek 22 AUGUST 2000 PHYSICS TODAY

$$\mathcal{L} = \frac{1}{4g^2} G_{\mu\nu}^a G_{\mu\nu}^a + \sum_j \bar{q}_j (i\gamma^\mu D_\mu + m_j) q_j$$

$$\text{where } G_{\mu\nu}^a \equiv \partial_\mu A_\nu^a - \partial_\nu A_\mu^a + if_{bc}^a A_\mu^b A_\nu^c$$

$$\text{and } D_\mu \equiv \partial_\mu + it^a A_\mu^a$$

That's it!

FIGURE 1. THE QCD LAGRANGIAN $\mathcal{L}$ displayed here is, in principle, a complete description of the strong interaction. But, in practice, it leads to equations that are notoriously hard to solve. Here m_j and q_j are the mass and quantum field of the quark of j th flavor, and A is the gluon field, with spacetime indices μ and ν and color indices a, b, c . The numerical coefficients f and t guarantee SU(3) color symmetry. Aside from the quark masses, the one coupling constant g is the only free parameter of the theory.

How many quarks? From PDG

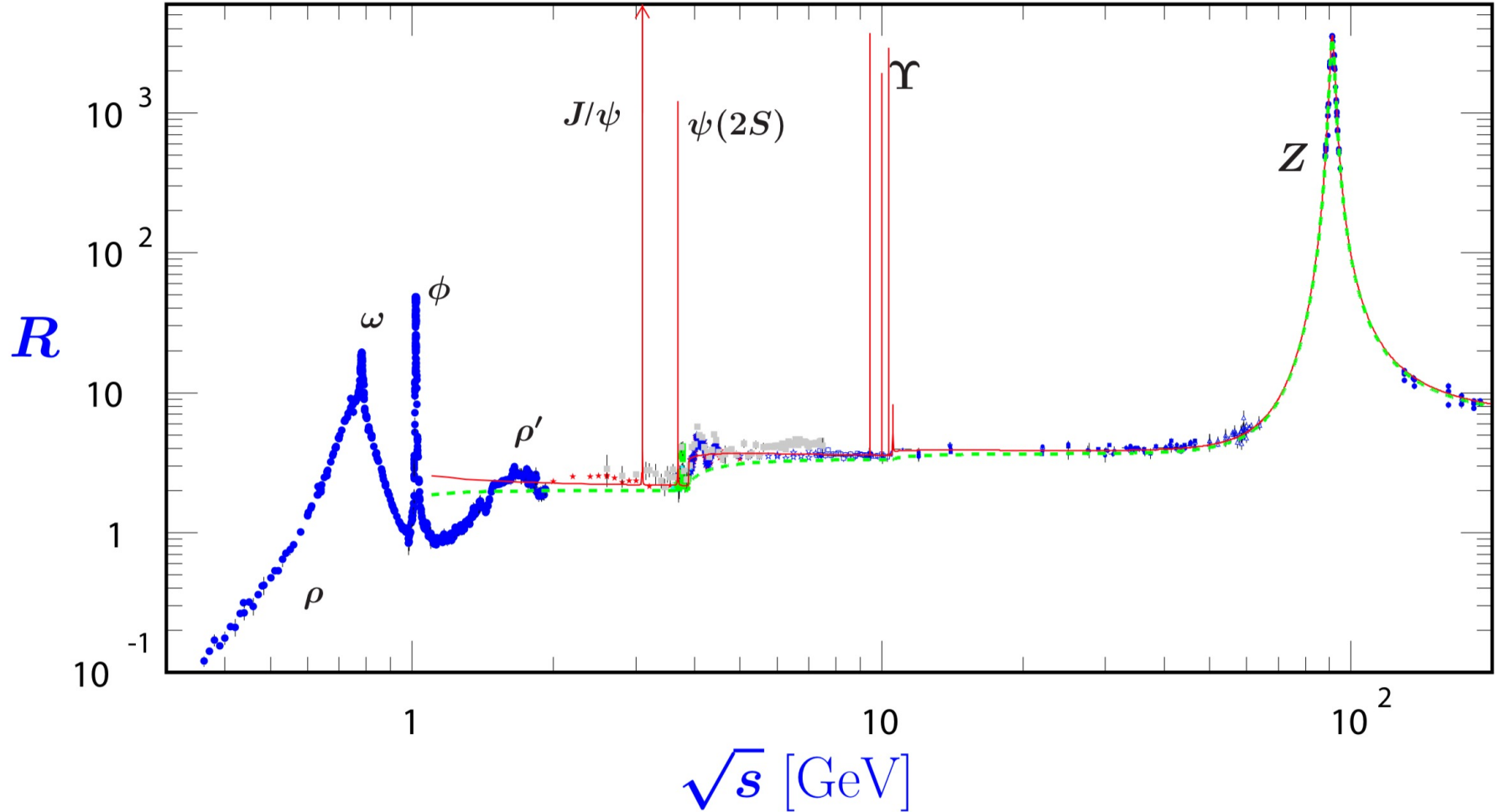
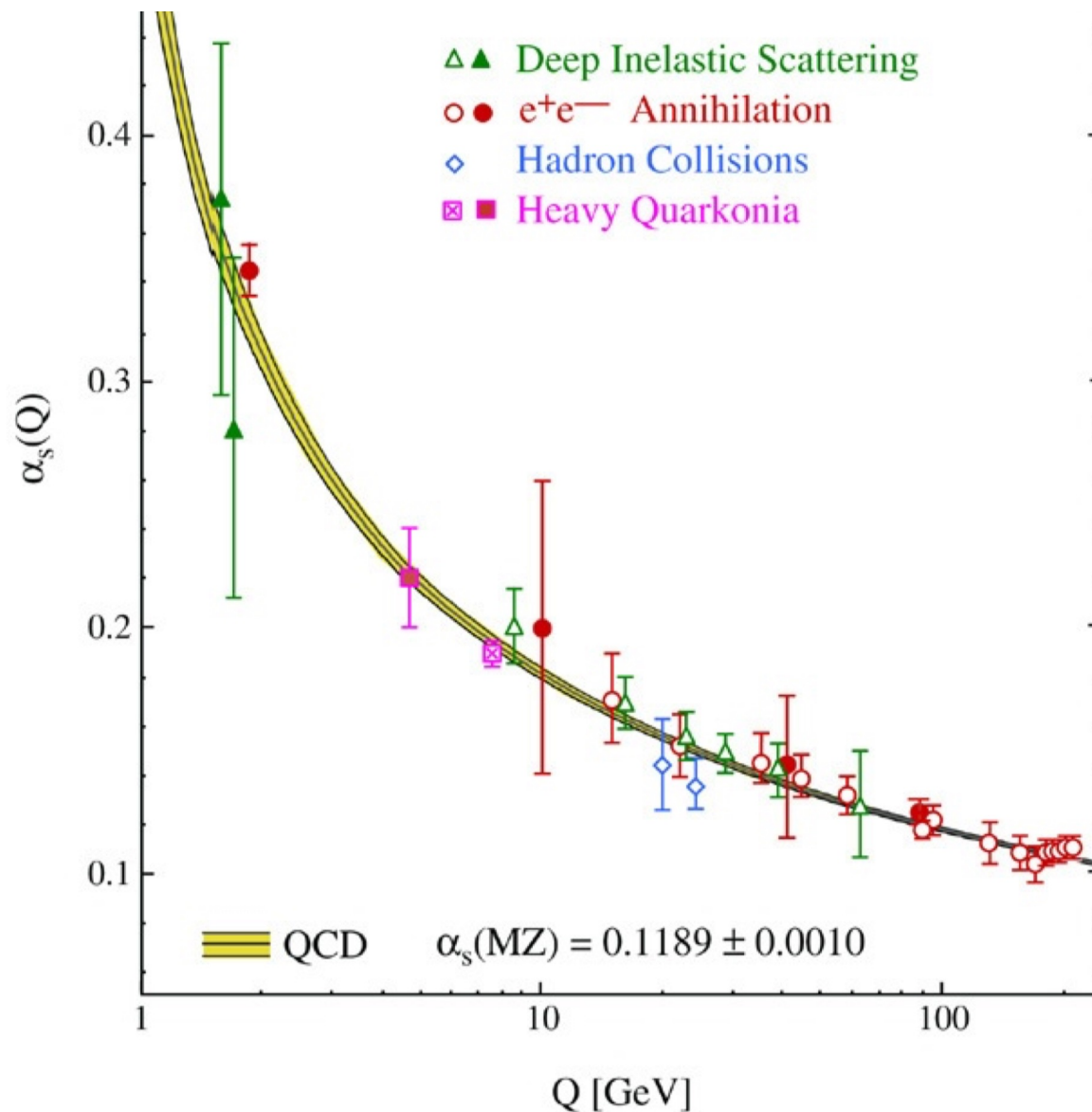


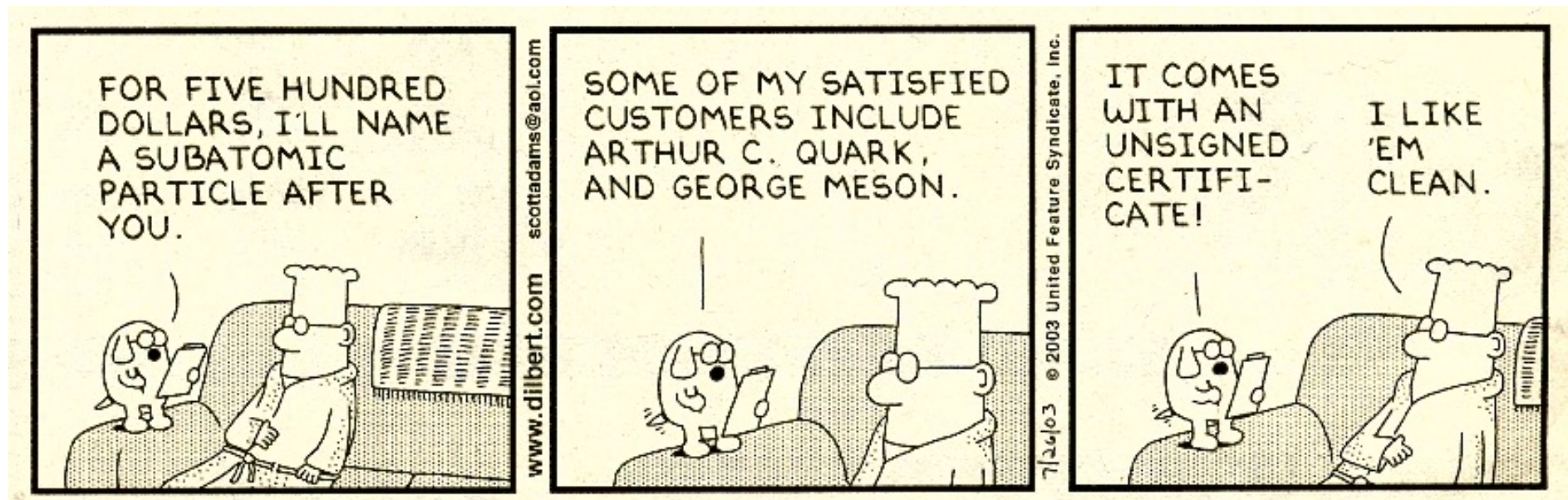
Figure 51.2: World data on the total cross section of $e^+e^- \rightarrow \text{hadrons}$ and the ratio $R(s) = \sigma(e^+e^- \rightarrow \text{hadrons}, s)/\sigma(e^+e^- \rightarrow \mu^+\mu^-, s)$. $\sigma(e^+e^- \rightarrow \text{hadrons}, s)$ is the experimental cross section corrected for initial state radiation and electron-positron vertex loops, $\sigma(e^+e^- \rightarrow \mu^+\mu^-, s) = 4\pi\alpha^2(s)/3s$. Data errors are total below 2 GeV and statistical above 2 GeV. The curves are an educative guide: the broken one (green) is a naive quark-parton model prediction, and the solid one (red) is 3-loop pQCD prediction (see “Quantum Chromodynamics” section of this *Review*, Eq. (9.7) or, for more details, K. G. Chetyrkin *et al.*, Nucl. Phys. **B586**, 56 (2000) (Erratum *ibid.* **B634**, 413 (2002))). Breit-Wigner parameterizations of J/ψ , $\psi(2S)$, and $\Upsilon(nS)$, $n = 1, 2, 3, 4$ are also shown. The full list of references to the original data and the details of the R ratio extraction from them can be found in [arXiv:hep-ph/0312114]. Corresponding computer-readable data files are available at <http://pdg.lbl.gov/current/xsect/>. (Courtesy of the COMPAS (Protvino) and HEPDATA (Durham) Groups, August 2017. Corrections by P. Janot (CERN) and M. Schmitt (Northwestern U.))

Running of the Strong Coupling Constant



$$\alpha_s(\mu) = \frac{12\pi}{(33 - 2n_f) \ln(\mu^2 / \Lambda^2)}$$

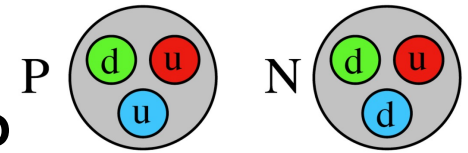
$$\Lambda_{QCD} \approx 0.25 \text{ GeV}$$



See also <http://particleadventure.org>

Hadronic Particle Zoo

- what can one build from quarks?



HADRONS							
Family Name	Particle Name	Particle Symbol	Antiparticle Symbol	Composition	Mass	Electric Charge	Lifetime in Seconds
baryon	proton	p or p ⁺	$\bar{p}$	uud	1,836	+1	stable
	neutron	n or n ⁰	$\bar{n}$	udd	1,839	0	887
	lambda	Λ^0	$\bar{\Lambda}^0$	uds	2,183	0	2.6×10^{-10}
	lambda-c	Λ_c^+	$\bar{\Lambda}_c^-$	udc	4,471	+1	2.1×10^{-13}
	lambda-b	Λ_b^0	$\bar{\Lambda}_b^0$	udb	11,000	0	1.1×10^{-12}
	sigma	Σ^+	$\bar{\Sigma}^-$	uus	2,328	+1	0.8×10^{-10}
		Σ^0	$\bar{\Sigma}^0$	$(ud \pm du)s$	2,334	0	7.4×10^{-20}
		Σ^-	$\bar{\Sigma}^+$	$\frac{\sqrt{2}}{2}dds$	2,343	-1	1.5×10^{-10}
	xi	Ξ^0	$\bar{\Xi}^0$	uss	2,573	0	2.9×10^{-10}
		Ξ^+	$\bar{\Xi}^-$	dss	2,585	-1	1.6×10^{-10}
	xi-c	Ξ_c^+	$\bar{\Xi}_c^-$	dsc	4,834	0	9.8×10^{-14}
		Ξ_c^0	$\bar{\Xi}_c^0$	usc	4,826	+1	3.5×10^{-13}
	omega	Ω^-	$\bar{\Omega}^+$	sss	3,272	-1	0.8×10^{-10}
	omega-c	Ω_c^0	$\bar{\Omega}_c^0$	ssc	5,292	0	6.4×10^{-14}
meson	pion	π^+	π^-	$u\bar{d}$	273	+1	2.6×10^{-8}
		π^0	π^0	$\frac{(u\bar{u} - d\bar{d})}{\sqrt{2}}$	264	0	8.4×10^{-17}
	kaon*	K^+	K^-	$u\bar{s}$	966	+1	1.2×10^{-8}
		K^0	$\bar{K}^0$	$d\bar{s}$	974	0	8.9×10^{-11}
	J/psi	J or Ψ	J or Ψ	$c\bar{c}$	6,060	0	1.0×10^{-20}
	omega	ω	ω	$\frac{(u\bar{u} + d\bar{d})}{\sqrt{2}}$	1,532	0	6.6×10^{-23}
	eta	η	η	$\frac{(u\bar{u} + d\bar{d})}{\sqrt{2}}$	1,071	0	3.5×10^{-22}
	eta-c	η_c	η_c	$c\bar{c}$	5,832	0	3.1×10^{-22}
	B	B^0	$\bar{B}^0$	$d\bar{b}$	10,331	0	1.6×10^{-12}
		B^+	B^-	$u\bar{b}$	10,331	+1	1.6×10^{-12}
	B-s	B_s^0	$\bar{B}_s^0$	$s\bar{b}$	10,507	0	1.6×10^{-12}
	D	D_0	$\bar{D}_0$	$c\bar{u}$	3,649	0	4.2×10^{-13}
		D^+	D^-	$c\bar{d}$	3,658	+1	1.1×10^{-12}
	D-s	D_s^+	D_s^-	$c\bar{s}$	3,852	+1	4.7×10^{-13}
	chi	χ_c^0	$\bar{\chi}_c^0$	$c\bar{c}$	6,687	0	3.0×10^{-22}
	psi	Ψ_c^0	$\bar{\Psi}_c^0$	$c\bar{c}$	7,213	0	1.5×10^{-20}
	upsilon	Y	Y	$b\bar{b}$	18,513	0	8.0×10^{-20}

*The neutral kaon is composed of two particles; the average lifetime of each particle is given.

Angular Momentum, Spin, Parity, ...

- Total angular momentum J is sum of orbital angular momentum L and Spin S
- Parity $P = +$ or $-$ = behavior under reflections
 - Antiparticles have opposite parity to particles
 - Angular momentum eigenstates with $L = 0, 1, 2, \dots$ have parity $+, -, +, \dots$
- ...and Charge Symmetry C

$\pi^\pm$

$$I^G(J^P) = 1^-(0^-)$$

Mass $m = 139.57039 \pm 0.00018$ MeV (S = 1.8)
 Mean life $\tau = (2.6033 \pm 0.0005) \times 10^{-8}$ s (S = 1.2)
 $c\tau = 7.8045$ m

 $\pi^\pm \rightarrow \ell^\pm \nu \gamma$ form factors ^[a]

$F_V = 0.0254 \pm 0.0017$
 $F_A = 0.0119 \pm 0.0001$
 F_V slope parameter $a = 0.10 \pm 0.06$
 $R = 0.059^{+0.009}_{-0.008}$

 π^- modes are charge conjugates of the modes below.

For decay limits to particles which are not established, see the section on
 Searches for Axions and Other Very Light Bosons.

 π^+ DECAY MODES

	Fraction (Γ_i/Γ)	Confidence level	p (MeV/c)
$\mu^+ \nu_\mu$	[b] (99.98770 $\pm$ 0.00004) %		30

 π^0

$$I^G(J^{PC}) = 1^-(0^{-+})$$

Mass $m = 134.9768 \pm 0.0005$ MeV (S = 1.1)
 $m_{\pi^\pm} - m_{\pi^0} = 4.5936 \pm 0.0005$ MeV
 Mean life $\tau = (8.43 \pm 0.13) \times 10^{-17}$ s (S = 1.2)
 $c\tau = 25.3$ nm

For decay limits to particles which are not established, see the appropriate
 Search sections (A^0 (axion) and Other Light Boson (X^0) Searches, etc.).

 π^0 DECAY MODES

	Fraction (Γ_i/Γ)	Scale factor/ Confidence level	p (MeV/c)
2γ	(98.823 $\pm$ 0.034) %	S=1.5	67
$e^+ e^- \gamma$	(1.174 $\pm$ 0.035) %	S=1.5	67

 $K^\pm$

$$I(J^P) = \frac{1}{2}(0^-)$$

Mass $m = 493.677 \pm 0.015$ MeV ^[o] (S = 2.8)
 Mean life $\tau = (1.2380 \pm 0.0020) \times 10^{-8}$ s (S = 1.8)
 $c\tau = 3.711$ m

 K^0

$$I(J^P) = \frac{1}{2}(0^-)$$

50% K_S , 50% K_L

Mass $m = 497.611 \pm 0.013$ MeV (S = 1.2)
 $m_{K^0} - m_{K^\pm} = 3.934 \pm 0.020$ MeV (S = 1.6)

 $D^\pm$

$$I(J^P) = \frac{1}{2}(0^-)$$

Mass $m = 1869.66 \pm 0.05$ MeV
 Mean life $\tau = (1033 \pm 5) \times 10^{-15}$ s
 $c\tau = 309.8$ μ m

 η

$$I^G(J^{PC}) = 0^+(0^{-+})$$

Mass $m = 547.862 \pm 0.017$ MeV
 Full width $\Gamma = 1.31 \pm 0.05$ keV

(... and B-mesons...)

$\rho(770)$

$$I^G(J^{PC}) = 1^+(1^- -)$$

See the review on "Spectroscopy of Light Meson Resonances."

T-Matrix Pole $\sqrt{s} = (761-765) - i(71-74) \text{ MeV}$

ρ^0 mass (Breit-Wigner) = $775.26 \pm 0.23 \text{ MeV}$ [g]

$\rho^\pm$ mass (Breit-Wigner) = $775.11 \pm 0.34 \text{ MeV}$ [h]

ρ^0 full width (Breit-Wigner) = $147.4 \pm 0.8 \text{ MeV}$ [g] (S = 2.0)

$\rho^\pm$ full width (Breit-Wigner) = $149.1 \pm 0.8 \text{ MeV}$ [h]

$\rho(770)$ DECAY MODES	Fraction (Γ_i/Γ)	Scale factor/ Confidence level	p (MeV/c)
$\pi\pi$	~ 100 %		363

$\phi(1020)$

$$I^G(J^{PC}) = 0^-(1^- -)$$

Mass $m = 1019.460 \pm 0.016 \text{ MeV}$

Full width $\Gamma = 4.249 \pm 0.013 \text{ MeV}$ (S = 1.1)

$\phi(1020)$ DECAY MODES	Fraction (Γ_i/Γ)	Scale factor/ Confidence level	p (MeV/c)
$K^+ K^-$	(49.9 $\pm$ 0.5) %	S=1.5	127
$K_L^0 K_S^0$	(33.6 $\pm$ 0.4) %	S=1.3	110
$\rho\pi + \pi^+\pi^-\pi^0$	(14.9 $\pm$ 0.4) %	S=1.3	—

$J/\psi(1S)$

$$I^G(J^{PC}) = 0^-(1^- -)$$

Mass $m = 3096.900 \pm 0.006 \text{ MeV}$

Full width $\Gamma = 92.6 \pm 1.7 \text{ keV}$ (S = 1.1)

(...+bottomium...)

...and all L>0 states...
(208 pages in PDG)

BARYONS:

Simple (Constituent) Quark Model

Flavor	Isospin I	I_3	Strangeness S	Charge Q	Baryon Number B
U	$1/2$	$+1/2$	0	$+2/3$	$1/3$
D	$1/2$	$-1/2$	0	$-1/3$	$1/3$
S	0	0	-1	$-1/3$	$1/3$

$$|\Delta^{++} \uparrow\rangle = |U \uparrow U \uparrow U \uparrow\rangle$$

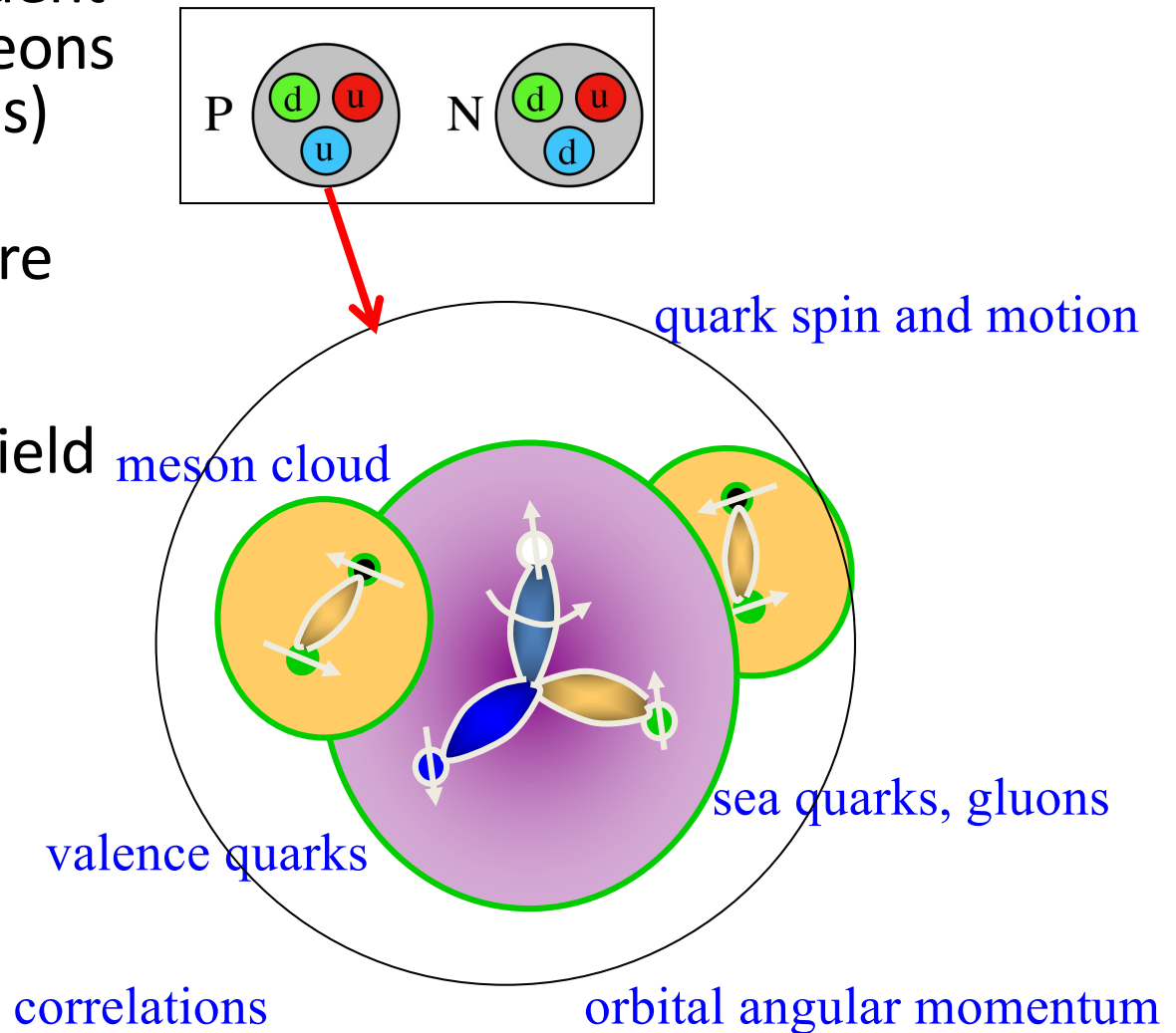
$$|\Delta^+ \uparrow\rangle = 1/\sqrt{3} (|U \uparrow U \uparrow D \uparrow\rangle + |U \uparrow D \uparrow U \uparrow\rangle + |D \uparrow U \uparrow U \uparrow\rangle)$$

The case of the proton is a bit more complicated, since the wave function cannot be symmetric in spin and flavor separately. The most intuitive way to derive the proton wave function is by observing that 2 of the 3 quarks are equal (U), and therefore their relative spin wave function should be symmetric also. This leads to the conclusion that the two U -quarks couple their spins to a total spin of one. Let's denote the case where this spin has a z-projection of +1 as $(UU \uparrow\uparrow) := |U \uparrow U \uparrow\rangle$, while the projection with $S_z = 0$ will be indicated by $(UU \Rightarrow) := 1/\sqrt{2} (|U \uparrow U \downarrow\rangle + |U \downarrow U \uparrow\rangle)$. We can now combine the spin 1/2 of the remaining D quark with the spin 1 of the UU pair in two ways to get total spin and projection 1/2; the proper way follows simply from insertion of the correct Clebsch-Gordon coefficients:

$$|P \uparrow\rangle = 1/\sqrt{3} \left(\sqrt{2} |(UU \uparrow\uparrow) D \downarrow\rangle - |(UU \Rightarrow) D \uparrow\rangle \right). \quad (2)$$

Hadron Structure

- Simple-most (constituent quark) model of nucleons (protons and neutrons)
- ... becomes much more complicated once we consider the full relativistic quantum field theory called QCD
- Effective theories: Quark model, χ PT, sum rules, ...
- and Lattice QCD!

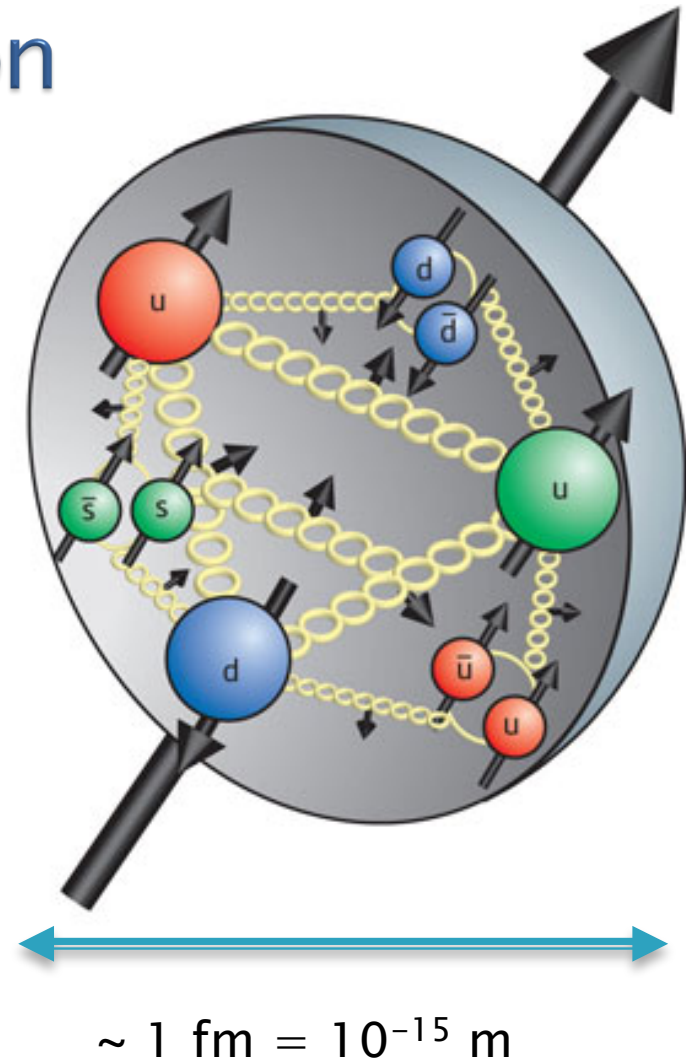


Reminder:

The structure of the nucleon

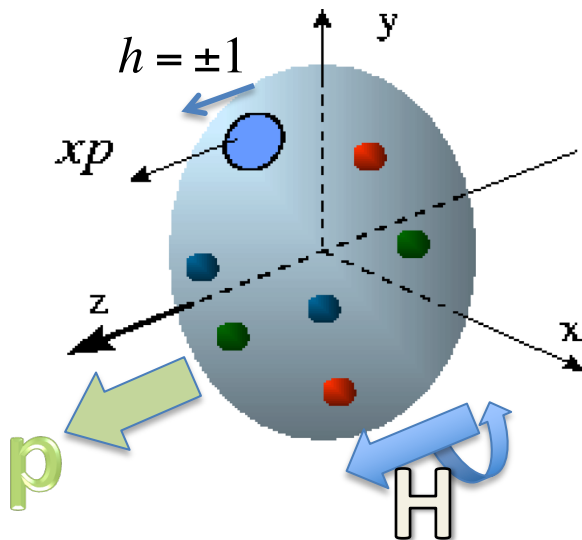
Three "valence" quarks plus gluons plus quark/anti-quark pairs.

- How do the quarks and gluons interact to form a proton or neutron?
- Where does the spin of the nucleon come from?
- Why is the radius of the proton different when measured by electrons and muons?
- Why are protons and neutrons modified in a nucleus?



Inclusive lepton scattering

Parton model: DIS can access



$$q(x; Q^2), \langle h \cdot H \rangle q(x; Q^2)$$

Traditional “1-D” Parton Distributions (PDFs) (integrated over many variables)

$$F_1(x) = \frac{1}{2} \sum_i e_i^2 q_i(x) \quad (\text{and } F_2(x) \approx 2xF_1(x)) \quad \text{Callan-Gross}$$

$$g_1(x) = \frac{1}{2} \sum_i e_i^2 \Delta q_i(x) \quad \left(\text{and } g_2(x) \approx -g_1(x) + \int_x^1 \frac{g_1(y)}{y} dy \right) \quad \text{Wandzura-Wilczek}$$

At finite Q^2 : pQCD evolution ($q(x, Q^2)$, $\Delta q(x, Q^2) \Rightarrow$ DGLAP equations), and gluon radiation

$$g_1(x, Q^2)_{pQCD} = \frac{1}{2} \sum_q^{N_f} e_q^2 [(\Delta q + \Delta \bar{q}) \otimes (1 + \frac{\alpha_s(Q^2)}{2\pi} \delta C_q) + \frac{\alpha_s(Q^2)}{2\pi} \Delta G \otimes \frac{\delta C_G}{N_f}]$$

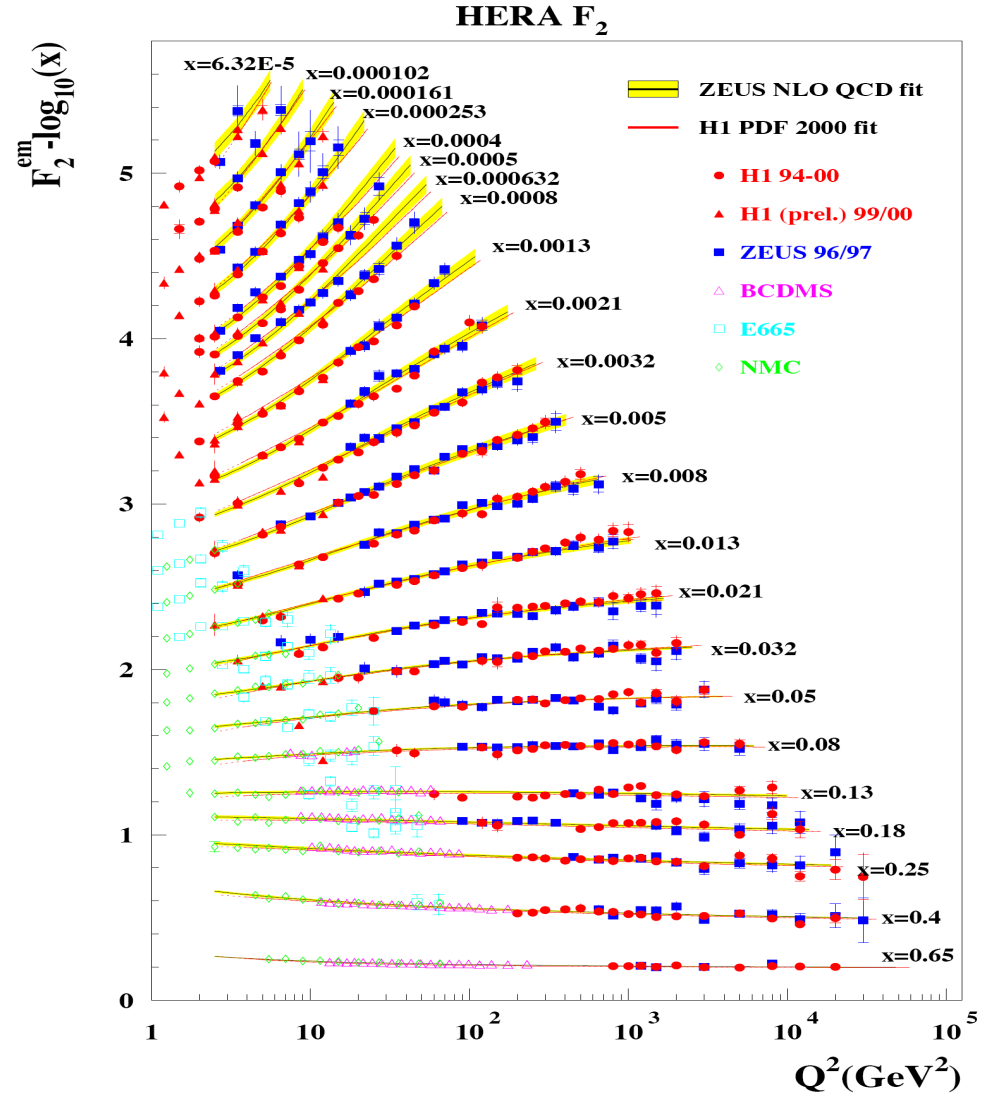
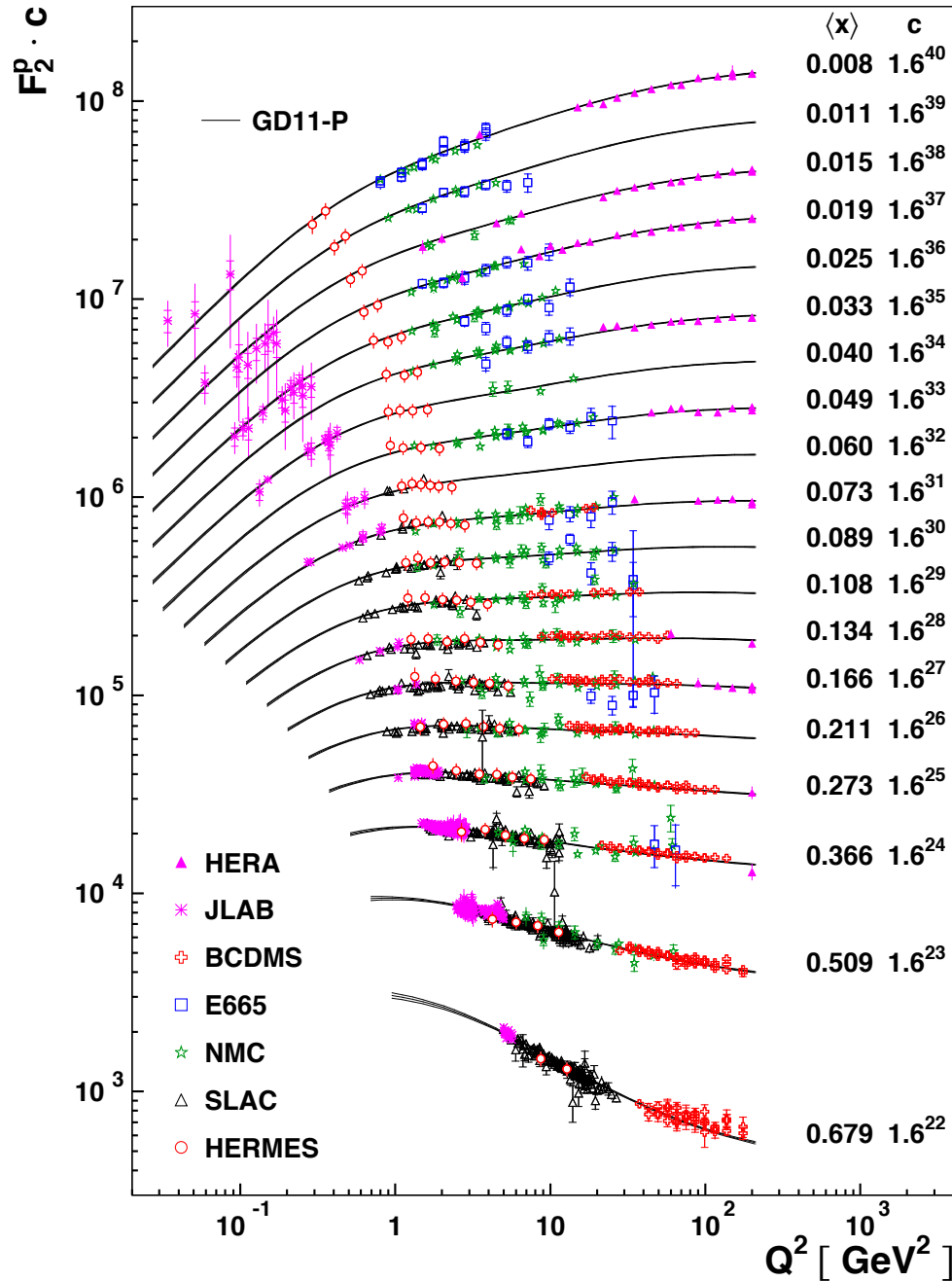
$\Rightarrow$ access to gluons. $\delta C_q, \delta C_G$ – Wilson coefficient functions

SIDIS: Tag the flavor of the struck quark with the leading FS hadron $\Rightarrow$ separate $q_i(x, Q^2)$, $\Delta q_i(x, Q^2)$

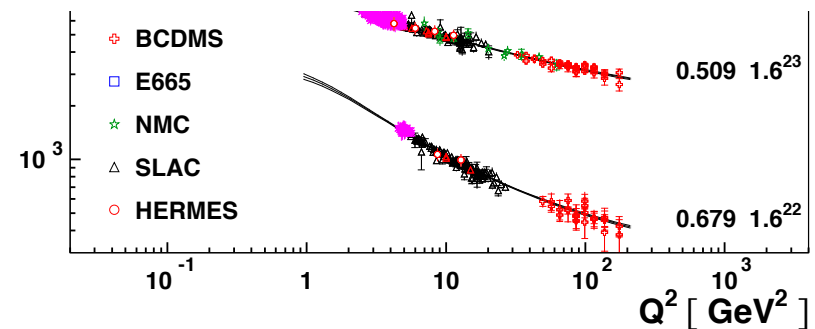
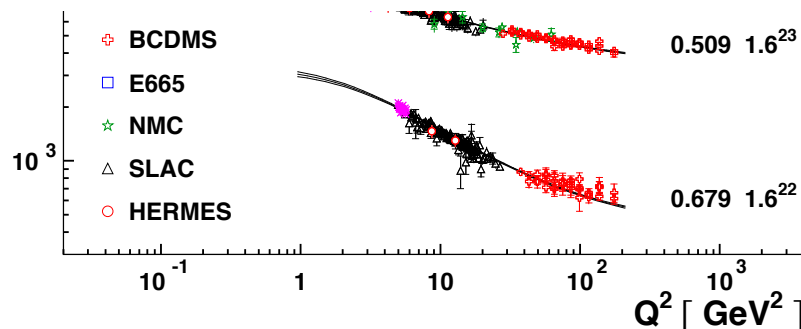
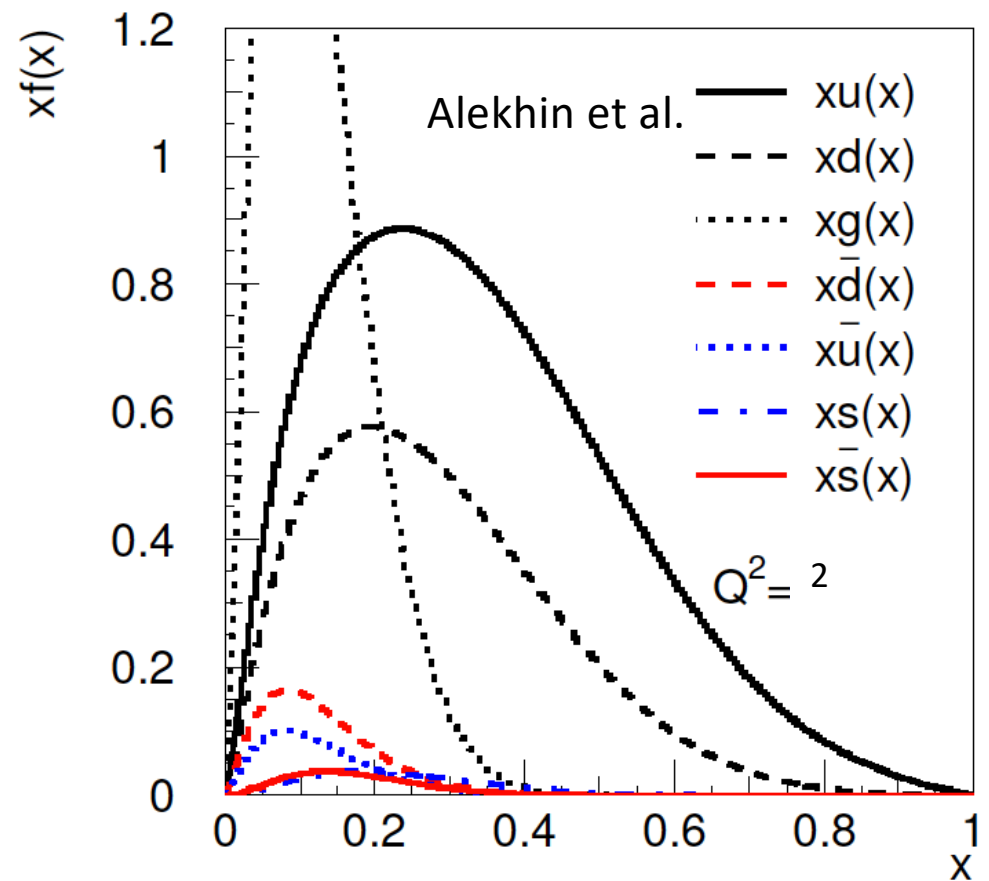
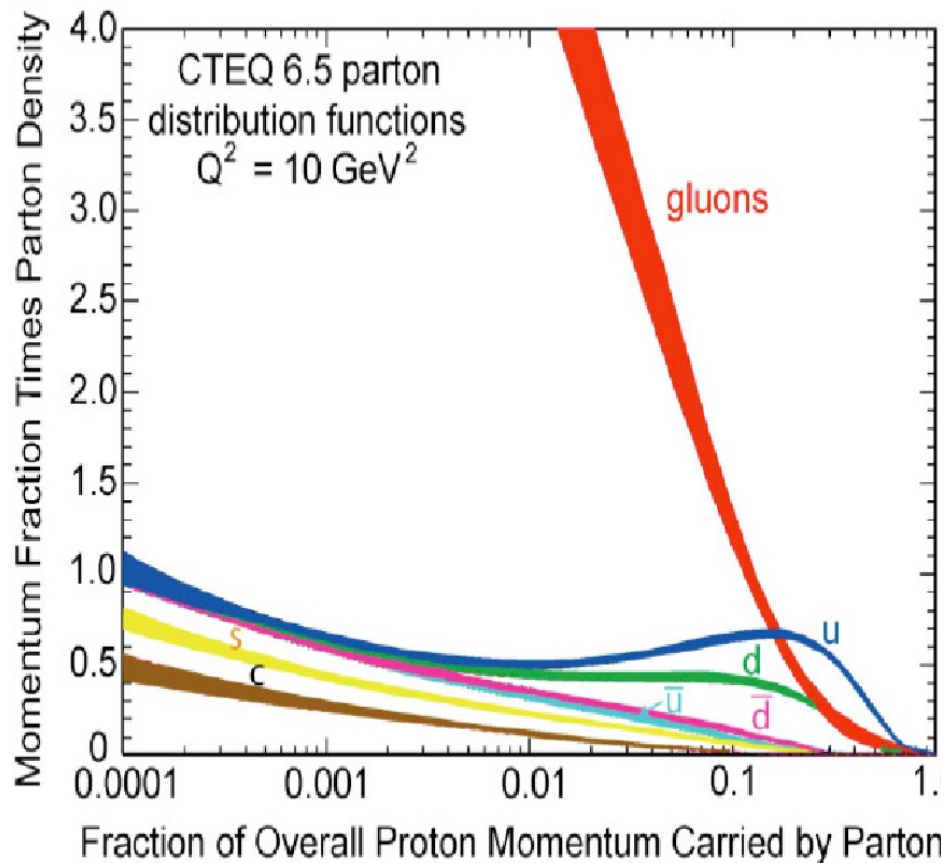
Jefferson Lab kinematics: $Q^2 \approx M^2 \Rightarrow$ target mass effects, higher twist contributions and resonance excitations

- Non-zero $R = \frac{F_2}{2xF_1} \left(\frac{4M^2x^2}{Q^2} + 1 \right) - 1$, $g_2^{HT}(x) = g_2(x) - g_2^{WW}(x)$
- Further Q^2 -dependence (power series in $\frac{1}{Q^n}$)

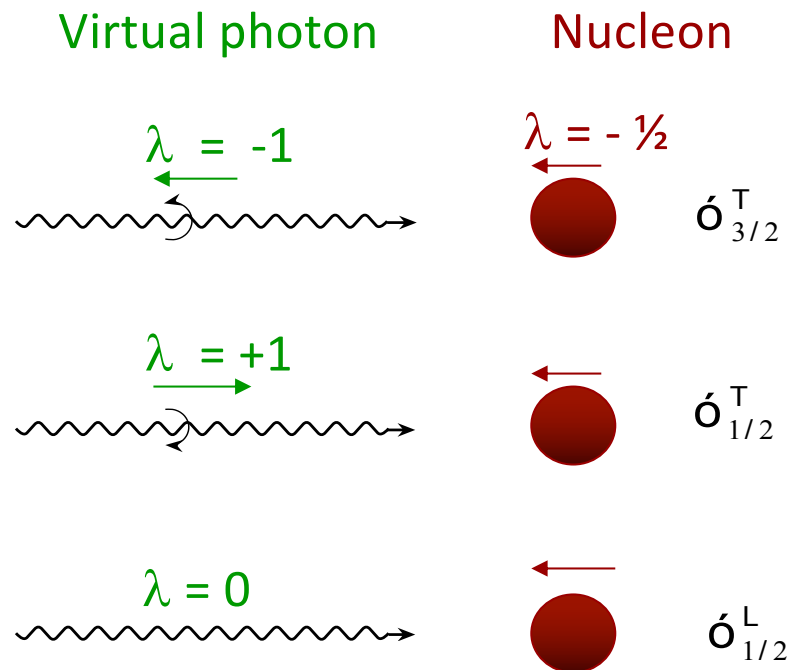
Results



Unpolarized Structure Functions



Virtual Photon Asymmetries



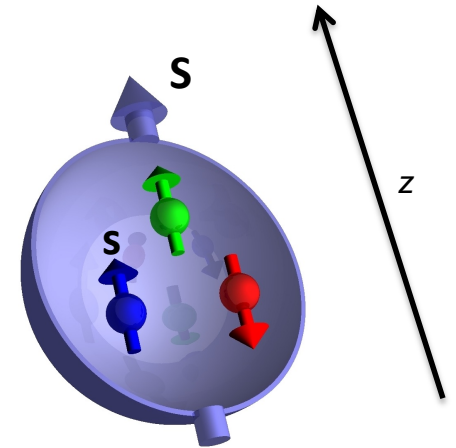
a) b)

$A_1 = \frac{\sigma_{1/2} - \sigma_{3/2}}{\sigma_T}$

$A_2 = \frac{\sigma_{LT'}}{\sigma_T}$

related to quark polarizations
 $\Delta q/q$

Polarized Parton Distribution Functions

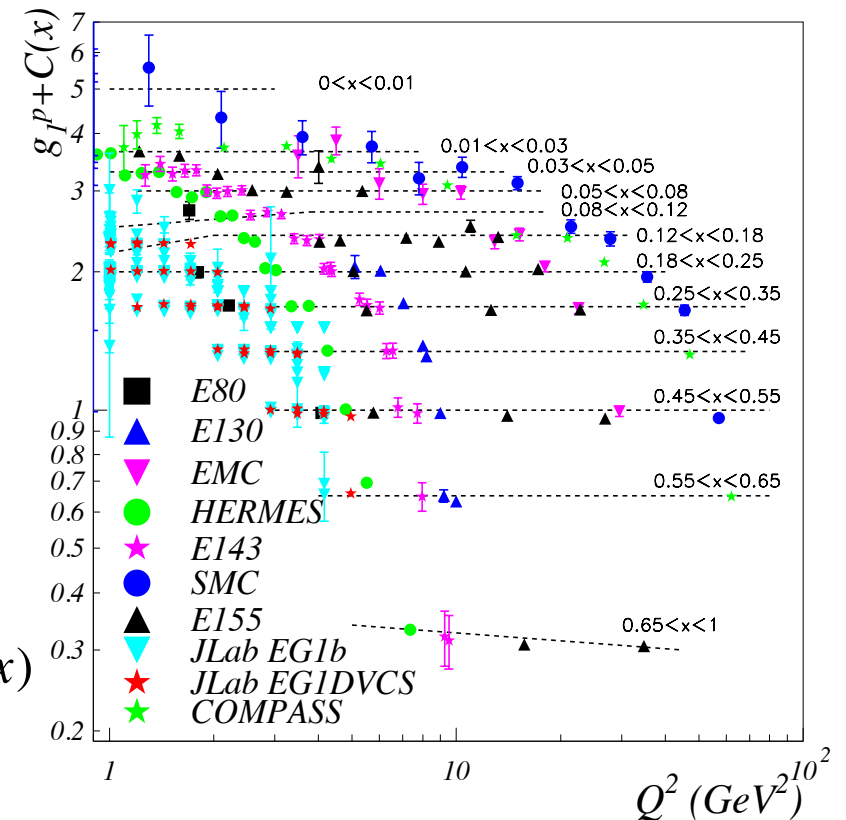


- Introduce two more quantities of interest:
 - Proton spin $\mathbf{S}$
 - Parton spin $\mathbf{s}$
 - Now we have 3 vectors: $\hat{z}, \vec{S}, \vec{s}$
 - **But:** Every observable must be a scalar
 - **And:** Spins are axial vectors!
 - **Finally:** Must treat longitudinal and transverse directions differently (boost)
 - 2 Pseudoscalars: $H = \vec{S} \cdot \hat{z}, h = \vec{s} \cdot \hat{z}$
 - 2 transverse (2D) axial vectors: $\vec{S}_\perp, \vec{s}_\perp$
 - 2nd Structure function

$$g_1^i(x) = \langle hH \rangle q_i(x) \text{ or } \langle hH \rangle G(x) = \Delta q_i(x) \text{ or } \Delta G(x)$$

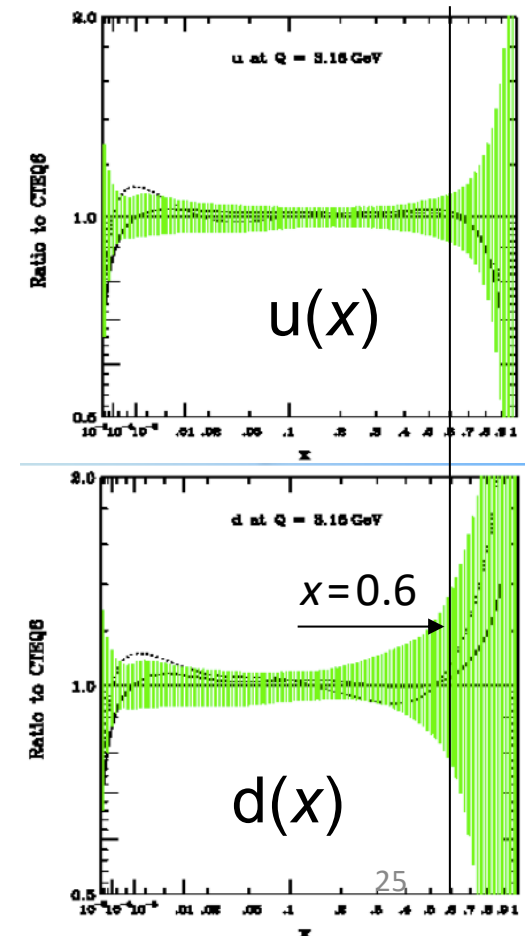
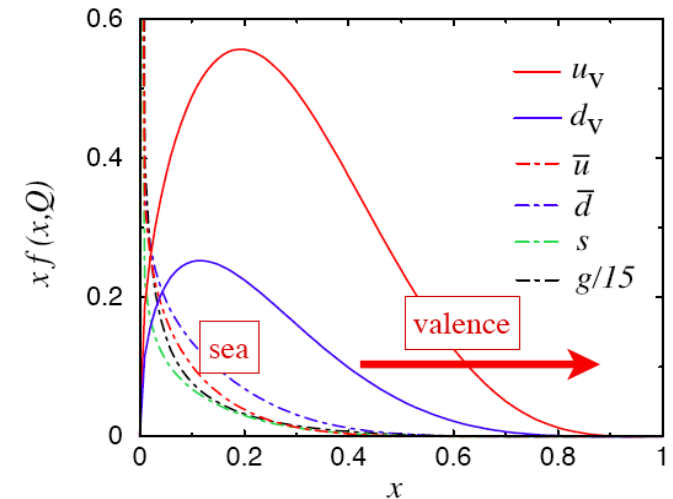
$$\Delta q_i = q_{\uparrow\uparrow}(x) - q_{\uparrow\downarrow}(x)$$

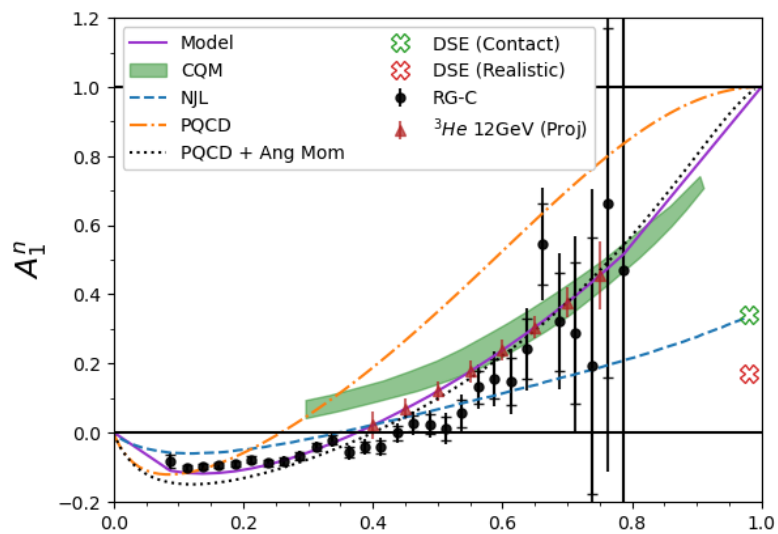
Can also form one more scalar: $T = \vec{S}_\perp \cdot \vec{s}_\perp$ (not measurable in DIS) $\rightarrow$ Transversity $h_1(x)$



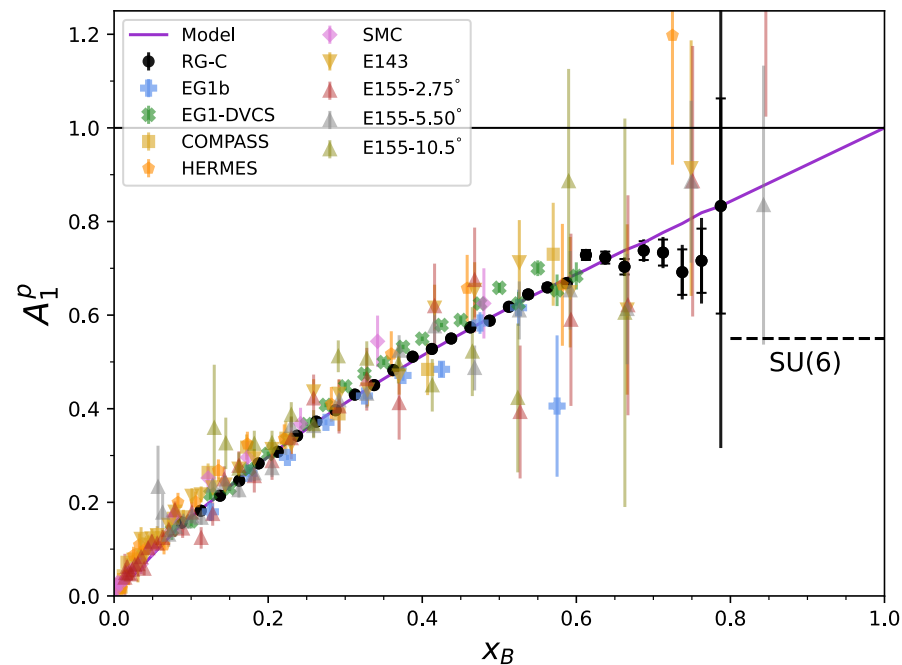
Valence PDFs

- Behavior of PDFs still unknown for $x \rightarrow 1$
 - SU(6): $d/u = 1/2$, $\Delta u/u = 2/3$, $\Delta d/d = -1/3$ for all x
 - Relativistic Quark model: Δu , Δd reduced
 - Hyperfine effect (1-gluon-exchange): Spectator spin 1 suppressed, $d/u \rightarrow 0$, $\Delta u/u \rightarrow 1$, $\Delta d/d \rightarrow -1/3$
 - Helicity conservation: $d/u \rightarrow 1/5$, $\Delta u/u \rightarrow 1$, $\Delta d/d \rightarrow 1$
 - Orbital angular momentum: can explain slower convergence to $\Delta d/d \rightarrow 1$
- Plenty of data on proton $\rightarrow$ mostly constraints on u and Δu
- Knowledge on d limited by lack of free neutron target (nuclear binding effects in d , ^3He)
- Large x requires very high luminosity and resolution; binding effects become dominant uncertainty for the neutron

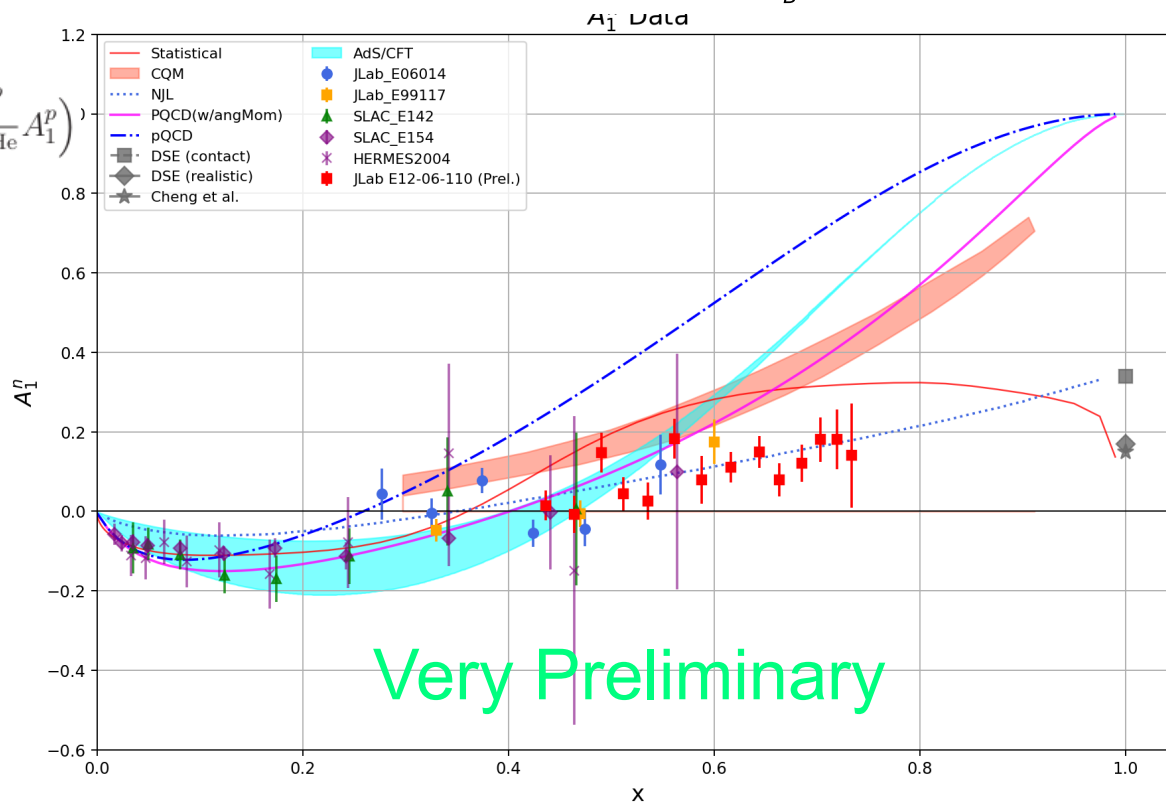




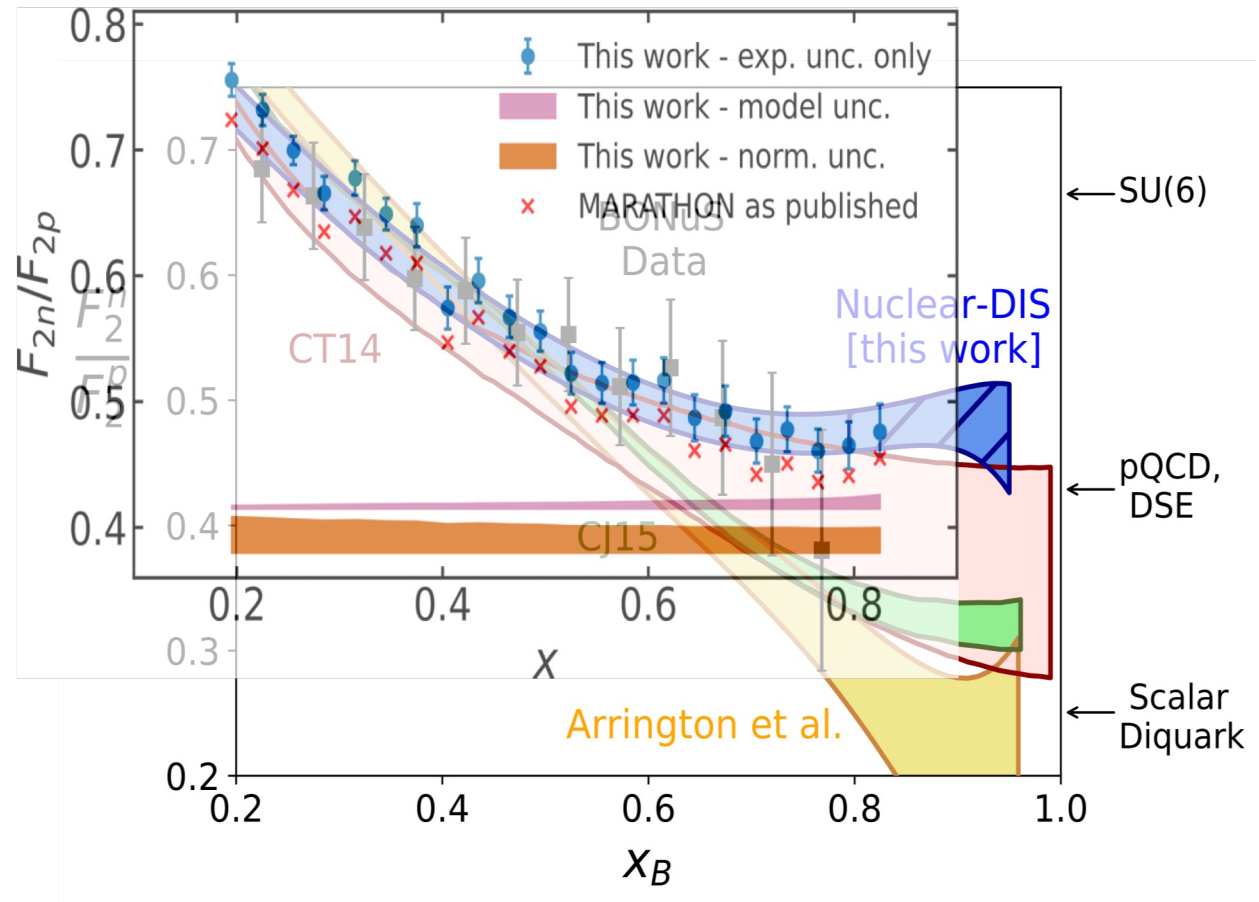
A_1^n extracted using effective polarizations
plus our own ^3He fit)



$$A_1^n = \frac{1}{P_n} \frac{F_2^{^3\text{He}}}{F_2^n} \left(A_1^{^3\text{He}} - 2P_p \frac{F_2^p}{F_2^{^3\text{He}}} A_1^p \right)$$



Very Preliminary



F2n/F2p based on F2d/F2p from JAM

