

**Nuclear Physics - Last Problem Set 10 – Solution**

**Problem 1)**

Since we are not supposed to count the strong coupling constant, there are only 2 constants describing electroweak coupling. Those could be  $g$  and  $g'$  or, more usefully,  $e$  and  $\sin\theta_w$ .

Secondly, we have to include some information about the Higgs sector – e.g. the mass of the observed Higgs boson and the vacuum expectation value of the Higgs field at the potential minimum. These in turn determine uniquely the masses of both  $W$ 's and  $Z$ 's (together with the couplings above).

The remaining constants are the Yukawa couplings of all 12 fundamental fermions to the Higgs field, i.e., their masses, and the 4 parameters each of the mixing matrix for the quark sector and for the neutrino sector. This would be 20 more parameters, for a total of 24 (25 if you include the strong coupling and 26 including gravity). Hardly a “fundamental theory”, but very successful. (NB: For completeness, we would also have to include a “QCD phase” which makes the total 27).

**Problem 2)**

Here are my answers:

${}^3\text{H} \rightarrow {}^3\text{He}$  (g.s.):  $1/2^+ \rightarrow 1/2^+$  transition; therefore  $\Delta l=0$ . Both Gamov-Teller and Fermi can (and DO) contribute. The two nuclei are isospin partners with the same wave functions (under exchange of n's with p's), therefore this transition is super-allowed.

${}^{20}\text{F} \rightarrow {}^{20}\text{Ne}$  (g.s.):  $2^+ \rightarrow 0^+$  transition; therefore  $\Delta l=2$ . This transition is twice forbidden. Both Gamov-Teller and Fermi contribute.

${}^{42}\text{Sc}$  (g.s.)  $\rightarrow$   ${}^{42}\text{Ca}$  (g.s.):  $0^+ \rightarrow 0^+$  transition; therefore  $\Delta l=0$ . This transition is allowed (and in fact even super-allowed) and pure Fermi.

$^{42}\text{Sc}$  ( $7^+$  isomer)  $\rightarrow$   $^{42}\text{Ca}$  (g.s.):  $7^+ \rightarrow 0^+$  transition; therefore  $\Delta l=6$ . This transition is six times forbidden and pure Gamov-Teller (no wonder it never happens).

$^{64}\text{Cu}$   $\rightarrow$   $^{64}\text{Zn}$  (g.s.):  $1^+ \rightarrow 0^+$  transition; therefore  $\Delta l=0$ . This transition is allowed and pure Gamov-Teller.

$^{84}\text{Br}$   $\rightarrow$   $^{84}\text{Kr}$  (g.s.):  $2^- \rightarrow 0^+$  transition; therefore  $\Delta l=1$ . This transition is once-forbidden and pure Gamov-Teller.

### **Problem 3)**

From Problem 2), we realize that the transition in question is super-allowed. Therefore, the statement on page 332 applies: the  $ft$  value of this transition is “roughly equal” to that of the free neutron decay. From equations 16.56,

16.59 and 16.60 we conclude that for the neutron  $f\tau = 0.47 \frac{\epsilon_0^5}{30} \cdot 880 \text{ s}$ . Now

“all” we have to do is figure out the  $f(E_0)$  for the transition in question. First note that in these formulas, all energies are total energies (rest mass plus kinetic energy). Therefore, in the case of the neutron,

$E_0 = m_n - m_p = 1.2903 \text{ MeV} \Rightarrow \epsilon_0 = 2.525 \Rightarrow f\tau = 1415 \text{ s}$ . To calculate  $f(E_0)$  for the transition  $^3\text{H} \rightarrow ^3\text{He}$  I use equation (16.55) (p. 278). Using the actual masses of the Helium-3 nucleus and the triton, I get  $E_0 = 0.5296 \Rightarrow \epsilon_0 = 1.0364$ . Obviously, all energies involved are only slightly bigger than the electron rest mass, and I can use the non-relativistic approximation

$$E_e = m_e c^2 + \frac{p_e^2}{2m_e} \Rightarrow dE_e = \frac{p_e}{m_e} dp_e. \text{ Therefore, the integral can be}$$

approximated as

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$$\begin{aligned}
 f(E_o) &= \frac{1}{m_e^5 c^{10}} \int_{m_e c^2}^{E_o} E_e c p_e (E_o - E_e)^2 dE_e = \\
 &= \frac{1}{m_e^5 c^9} \int_0^{\sqrt{2m_e(E_o - m_e c^2)}} \left( m_e c^2 + \frac{p_e^2}{2m_e} \right) \frac{p_e^2}{m_e} \left( E_o - m_e c^2 - \frac{p_e^2}{2m_e} \right)^2 dp_e = \\
 &= \frac{1}{m_e^5 c^9} \int_0^{0.1379 \text{ MeV}/c} \left( \frac{p_e^8}{8m_e^4} + [m_e c^2 - 2(E_o - m_e c^2)] \frac{p_e^6}{4m_e^3} + \right. \\
 &\quad \left. \left[ (E_o - m_e c^2)^2 - 2m_e c^2 (E_o - m_e c^2) \right] \frac{p_e^4}{2m_e^2} + m_e c^2 (E_o - m_e c^2)^2 \frac{p_e^2}{m_e} \right) dp_e \\
 &= \frac{1}{m_e^5 c^9} \left[ \frac{p_e^9}{72m_e^4} + 0.4738 \text{ MeV} \frac{p_e^7}{28m_e^3} - 0.0187 \text{ MeV}^2 \frac{p_e^5}{10m_e^2} + 0.0001768 \text{ MeV}^3 \frac{p_e^3}{3m_e} \right]_0^{0.1379 \text{ MeV}/c} \\
 &= 2.01 \cdot 10^{-6}
 \end{aligned}$$

This leads to a lifetime of  $\tau=22.3$  years or a half-life of 15.5 years. This is not too far off the measured half-life of 12.3 years.